

Lecture 4 – Mechanical work and energy

Purdue ME 200, Thermodynamics I

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Outline

Differentials

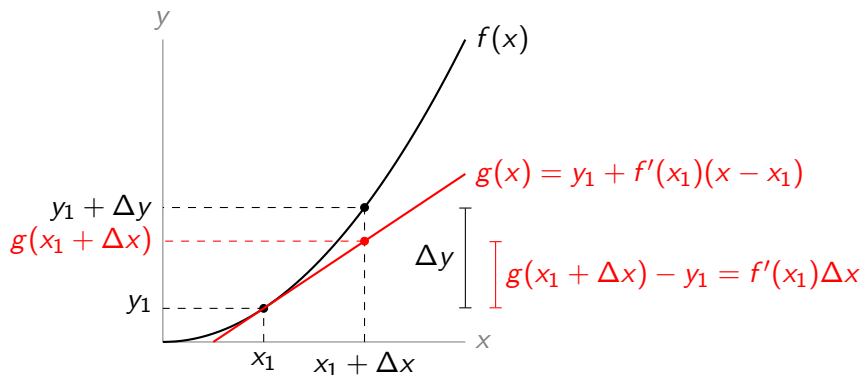
Mechanical work

Conservation of mechanical energy

What are differentials?

- a **differential** dx is just a tiny (infinitesimal) change in x
- dx can be viewed as the limit of a finite Δx as $\Delta x \rightarrow 0$
- differentials often show up under integral signs
(to calculate total change, add up a bunch of tiny changes)

Differentials in one dimension



- suppose $y = f(x)$ and g is the linear approximation to f at x_1
- when x changes by Δx , y changes by a true amount Δy
- for small Δx , the true $\Delta y \approx f'(x_1)\Delta x$

Differentials in one dimension (continued)

- as $\Delta x \rightarrow 0$, the approximation $\Delta y \approx f'(x_1)\Delta x$ becomes exact
- at any arbitrary x , we write

$$dy = f'(x)dx \text{ or } dy = \left(\frac{dy}{dx}\right) dx$$

- from the fundamental theorem of calculus,

$$\int_{y_1}^{y_2} dy = \int_{x_1}^{x_2} f'(x)dx = f(x_2) - f(x_1) = y_2 - y_1$$

- more compactly (but less precisely), we can write

$$\int dy = \Delta y$$

Differentials in two dimensions

- consider a variable $z = f(x, y)$
- the differential of z is

$$dz = \left(\frac{\partial f}{\partial x} \right)_y dx + \left(\frac{\partial f}{\partial y} \right)_x dy$$

- in 2D, $\int dz = \Delta z$ only holds for **exact** differentials
- only properties (p , v , T , ...) have exact differentials
- other things (work W , heat Q) have **inexact** differentials
- we write inexact differentials with δ instead of d
- we never write $\int \delta W = \Delta W$ or $\int \delta Q = \Delta Q$

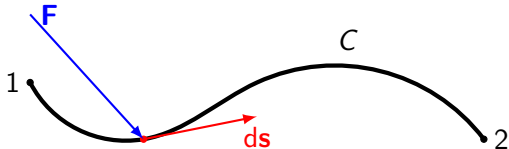
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General form of mechanical work



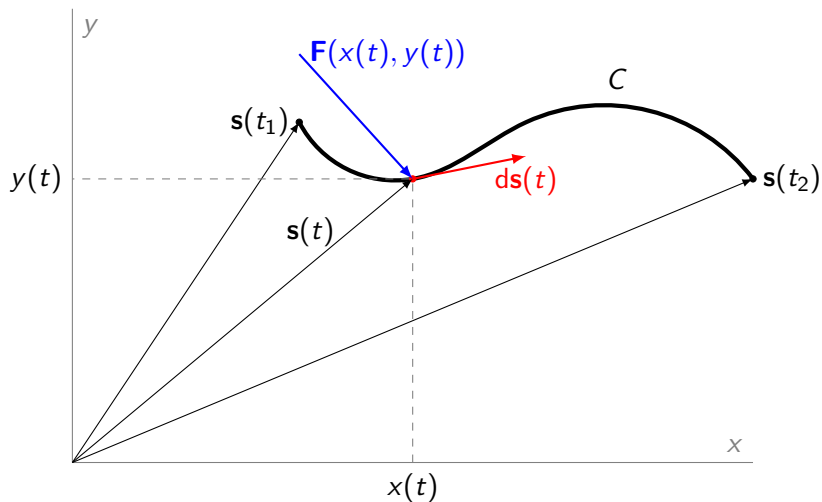
- work done by force \mathbf{F} on body moved over path C is

$$W = \int_C \mathbf{F} \cdot d\mathbf{s}$$

- differential displacement $d\mathbf{s}$ always points along path
- dot product picks out the component of \mathbf{F} along path
- if path starts at point 1 and ends at point 2, we may write

$$W = \int_1^2 \mathbf{F} \cdot d\mathbf{s}$$

Parameterization in terms of t



$$W = \int_{t_1}^{t_2} \mathbf{F}(x(t), y(t)) \cdot ds(t)$$

Calculating mechanical work

- force can be written as

$$\mathbf{F}(x(t), y(t)) = F_x(x(t), y(t))\mathbf{i} + F_y(x(t), y(t))\mathbf{j}$$

- path C contains all position vectors $\mathbf{s}(t)$ over $t_1 \leq t \leq t_2$
- position vector can be written as $\mathbf{s}(t) = x(t)\mathbf{i} + y(t)\mathbf{j}$
- so differential displacement along path is

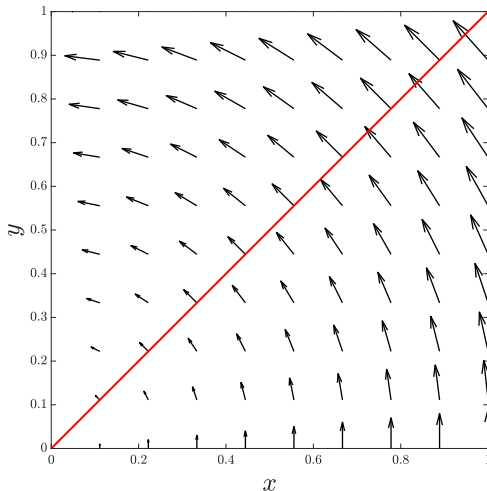
$$d\mathbf{s}(t) = dx(t)\mathbf{i} + dy(t)\mathbf{j} = x'(t)dt\mathbf{i} + y'(t)dt\mathbf{j}$$

- and work done by force \mathbf{F} over path C is

$$W = \int_{t_1}^{t_2} [F_x(x(t), y(t))x'(t) + F_y(x(t), y(t))y'(t)] dt$$

Example

- suppose force is $\mathbf{F}(x(t), y(t)) = -y(t)\mathbf{i} + x(t)\mathbf{j}$
- and path is $\mathbf{s}(t) = t\mathbf{i} + t\mathbf{j}$, $0 \leq t \leq 1$



Solution

- force is perpendicular to path, so $\mathbf{F} \cdot d\mathbf{s} = 0$ and $W = 0$
- does the math agree?

$$x(t) = t, \quad y(t) = t$$

$$x'(t) = 1, \quad y'(t) = 1$$

$$F_x(x(t), y(t)) = -y(t) = -t$$

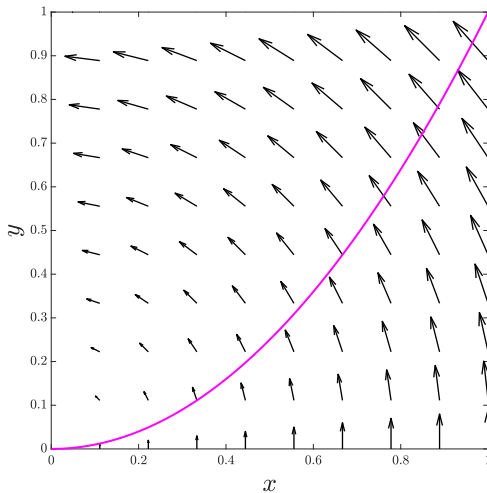
$$F_y(x(t), y(t)) = x(t) = t$$

- therefore,

$$\begin{aligned} W &= \int_{t_1}^{t_2} [F_x(x(t), y(t))x'(t) + F_y(x(t), y(t))y'(t)] dt \\ &= \int_0^1 [(-t)(1) + (t)(1)] dt = 0 \end{aligned}$$

Modified example

- what if **path** instead is $\mathbf{s}(t) = t\mathbf{i} + t^2\mathbf{j}$, $0 \leq t \leq 1$



Modified solution

- force is no longer perpendicular to path, so $\mathbf{F} \cdot d\mathbf{s} \neq 0$
- along the modified path,

$$x(t) = t, \quad y(t) = t^2$$

$$x'(t) = 1, \quad y'(t) = 2t$$

$$F_x(x(t), y(t)) = -y(t) = -t^2$$

$$F_y(x(t), y(t)) = x(t) = t$$

- so the work is

$$\begin{aligned} W &= \int_{t_1}^{t_2} [F_x(x(t), y(t))x'(t) + F_y(x(t), y(t))y'(t)] dt \\ &= \int_0^1 [(-t^2)(1) + (t)(2t)] dt \\ &= \int_0^1 (2t^2 - t^2) dt = \frac{1}{3} \neq 0 \end{aligned}$$

Work is path-dependent

- these examples show that work is **path-dependent**
- we can't know the work done without knowing the path taken
- this implies that **work is not a property**
- as a reminder, we write differential work as δW , not dW
- and we never write $\int \delta W = \Delta W$

Work and power

- power \dot{W} is defined as the time derivative of work
- from the fundamental theorem of calculus, $W = \int_{t_1}^{t_2} \dot{W} dt$
- earlier, we found that

$$W = \int_{t_1}^{t_2} [F_x(x(t), y(t))x'(t) + F_y(x(t), y(t))y'(t)] dt$$

- the integrand is the dot product of $\mathbf{F}(x(t), y(t))$ and velocity,

$$\mathbf{v}(t) = x'(t)\mathbf{i} + y'(t)\mathbf{j}$$

- so $W = \int_{t_1}^{t_2} \mathbf{F} \cdot \mathbf{v} dt$, and therefore

$$\dot{W} = \mathbf{F} \cdot \mathbf{v}$$

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Differentials

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Mechanical work and energy in one dimension

- consider a body of mass m at height y and (vertical) speed v
- its kinetic energy (KE) is $\frac{1}{2}mv^2$
- its gravitational potential energy (PE) is mgy
(defining PE = 0 at $y = 0$)
- suppose some other forces, summing to F , act on the body
- if the body moves from y_1 to y_2 , those forces do work

$$W = \int_{y_1}^{y_2} F dy$$

Conservation of energy in one dimension



- Newton's second law: $F - mg = ma$
- integrate both sides over y from y_1 to y_2 :

$$\int_{y_1}^{y_2} (F - mg) dy = m \int_{y_1}^{y_2} a dy$$

- the left-hand side is

$$\int_{y_1}^{y_2} (F - mg) dy = \int_{y_1}^{y_2} F dy - mg(y_2 - y_1) = W - \Delta PE$$

Conservation of energy in one dimension (continued)

- the right-hand side is

$$m \int_{y_1}^{y_2} a dy = m \int_{y_1}^{y_2} \frac{dv}{dt} dy = m \int_{y_1}^{y_2} \frac{dv}{dy} \frac{dy}{dt} dy = m \int_{y_1}^{y_2} v \frac{dv}{dy} dy$$

- change variables of integration from y to v :
 - ◇ $v = v_1$ at $y = y_1$, and $v = v_2$ at $y = y_2$
 - ◇ since v is a function of y , its differential is $dv = (dv/dy)dy$
- then the integral becomes

$$m \int_{v_1}^{v_2} v dv = m \left[\frac{1}{2} v^2 \right]_{v_1}^{v_2} = \frac{1}{2} m (v_2^2 - v_1^2) = \Delta KE$$

Conservation of energy in one dimension (continued)

- putting the left- and right-hand sides together gives

$$\Delta KE + \Delta PE = W$$

- this is conservation of (mechanical) energy in one dimension
- although the derivation was 1D, the result holds in 2D and 3D
- this is a special case of the **1st law** of thermodynamics
- we'll progressively generalize the 1st law to include
 - ◇ other types of work
 - ◇ internal energy
 - ◇ heat transfer
 - ◇ mass transfer

Sign convention

- in this derivation, W was work done by forces on body
- we usually take W to be work done by system on surroundings (historically, thermo focused on work done by heat engines)
- with this sign convention, 1st law becomes

$$\Delta KE + \Delta PE = -W$$

Example

1,000 kg of water, initially at rest, flows through a vertical displacement of 60 m in a frictionless pipe. At the bottom, the water flows through a turbine, which it exits at 20 m/s.

- (a) How much mechanical work does the water do on the turbine?
- (b) If this takes 10 s, what is the average mechanical power?

Given and find

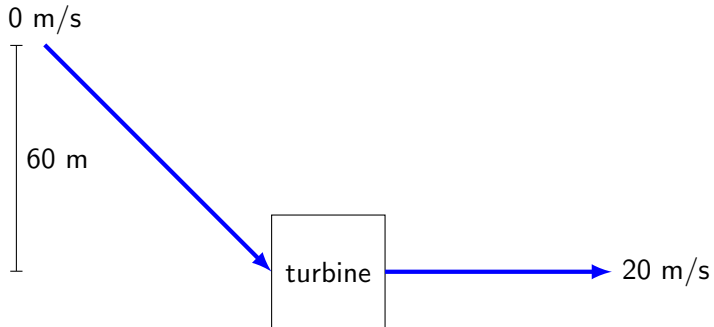
- **given**

- ◇ $m = 1000 \text{ kg}$
- ◇ $\Delta z = 60 \text{ m}$
- ◇ $v_1 = 0 \text{ m/s}$
- ◇ $v_2 = 20 \text{ m/s}$

- **find**

- (a) mechanical work W done by water on turbine
- (b) average mechanical power \dot{W}_{av} over $\Delta t = 10 \text{ s}$

System diagram



Assumptions and basic equations

- **assumptions**

- ◇ frictionless pipe (so no energy lost inside)

- **basic equations**

- ◇ $\Delta KE + \Delta PE = -W$

- ◇ also definitions: $KE = \frac{1}{2}mv^2$, $PE = mgz$

Solution

- change in kinetic energy is

$$\Delta KE = \frac{1}{2}m(v_2^2 - v_1^2) = \frac{1}{2}(1000\text{kg})[(20\text{m/s})^2 - (0\text{m/s})^2] = 200\text{kJ}$$

- change in potential energy is

$$\Delta PE = mg\Delta z = (1000\text{kg})(10\text{m/s}^2)(0\text{m} - 60\text{m}) = -600\text{kJ}$$

- mechanical work done by water on turbine is

$$W = -(\Delta KE + \Delta PE) = 400\text{kJ}$$

- average mechanical power is

$$\dot{W}_{\text{av}} = \frac{W}{\Delta t} = \frac{400\text{kJ}}{10\text{s}} = 40\text{kW}$$